

(1) 田路

(2) 全エネルギーが \$E\$ 以下の状態数

$$(2.1) \quad \Omega_0(E) = \frac{V^N}{N!} \int_{H \leq E} \frac{d^3 p_1 \dots d^3 p_N}{h^{3N}} = \frac{V^N}{N! h^{3N}} \left(\sqrt{2mE} \right)^{3N} \frac{\pi^{\frac{3N}{2}}}{\Gamma(\frac{3N}{2} + 1)} \dots (1)$$

$$H(p_1, \dots, p_N) = \frac{1}{2m} \sum_{i=1}^N p_i^2$$

エントロピー \$S(E, U)\$ は、状態数 \$\Omega\$ に対して

$$S(E, U) \approx k_B \log \Omega_0(E) \dots (2)$$

ここで、\$N \gg 1\$ のとき、\$S(E, U)\$ は、

$$S(E, U) = k_B N \log \frac{V}{U} + \frac{3}{2} k_B N \log \left(\frac{4\pi m}{3} \frac{E}{N} \right) + \frac{5}{2} k_B T \dots (3)$$

(2.2) 熱力学第一法則より \$dU = T dS - P dV\$ より

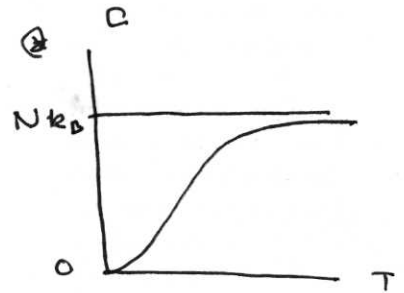
$$\frac{P}{T} = \left(\frac{\partial S}{\partial U} \right)_V = \frac{k_B N}{U} \quad E \text{ の関数として、} \dots$$

$$PV = k_B N T \quad E \text{ の関数として}$$

(2.3) 同様に

$$\frac{1}{T} = \left(\frac{\partial S}{\partial U} \right)_V = \frac{3}{2} k_B N \frac{1}{U}$$

$$(2.5) \quad U = \frac{3}{2} k_B N T$$



(3) (3.1) $Z = (1 + e^{-\beta \epsilon})^N$

(3.2) $U = - \frac{\partial \log Z}{\partial \beta} = \frac{\epsilon e^{-\beta \epsilon} N}{1 + e^{-\beta \epsilon}} = \frac{\epsilon N}{e^{\beta \epsilon} + 1}$

(3.3) $F = -k_B T \log Z = -k_B T N \log (1 + e^{-\beta \epsilon})$
 $S = - \frac{\partial F}{\partial T} = k_B N \log (1 + e^{-\beta \epsilon}) + k_B T N \frac{\epsilon}{k_B T^2} \frac{e^{-\beta \epsilon}}{1 + e^{-\beta \epsilon}}$

$$= k_B N \log (1 + e^{-\beta \epsilon}) + \frac{\epsilon N}{T} \frac{1}{e^{\beta \epsilon} + 1}$$

(4) (4.1) $Z = \sum_{n_1, \dots, n_N} e^{-\beta \hbar \omega (n_1 + \dots + n_N + \frac{N}{2})} = \left(\sum_{n=0}^{\infty} e^{-\beta \hbar \omega (n + \frac{1}{2})} \right)^N$
 $= \left(\frac{1}{2 \operatorname{sh} \frac{\beta \hbar \omega}{2}} \right)^N$

(4.2) $U = - \frac{\partial}{\partial \beta} \log Z = N \frac{\partial}{\partial \beta} \left(2 \operatorname{sh} \frac{\beta \hbar \omega}{2} \right) = \frac{N \hbar \omega}{2} \operatorname{coth} \frac{\beta \hbar \omega}{2}$

(4.3) $C = \frac{\partial U}{\partial T} = \frac{N}{k_B T^2} \left(\frac{\hbar \omega}{2} \right)^2 \frac{1}{\operatorname{sh}^2 \frac{\beta \hbar \omega}{2}}$